✴✴ DIRECT FORMULA FOR $\bar{E}$ in H·O ✴✴

$$\bar{E} = \frac{1}{2}kT \quad \text{for each term} \ (K\cdot E \& P\cdot E)$$

e.g. In 3-D, 3 K·E & 3 P·E term

$$\bar{E} = kT + kT + kT = 3kT$$

$$\boxed{\bar{E} = 3kT}$$

If particles are non interacting

$V=0$

$$\Rightarrow \bar{E} = \frac{1}{2}kT + \frac{1}{2}kT + \frac{1}{2}kT$$

$$\Rightarrow \boxed{\bar{E} = \frac{3}{2}kT}$$

For H.O.

$$Z = \sum_{n=0}^{\infty} e^{-\beta(n+\frac{1}{2})\hbar\omega}$$

$$Z = e^{-\hbar\omega\beta/2} + e^{-3\hbar\beta/2} + \cdots$$

G·P series

$$a = e^{-\hbar\omega\beta/2}, \quad r = e^{-\beta\hbar\omega}$$

$$\therefore z = \frac{e^{-\beta\hbar\omega/2}}{1-e^{-\beta\hbar\omega}}$$

For 3D H·O, $z_N = \frac{1}{LN}\left(\frac{kT}{\hbar\omega}\right)^{3N}$

Average value of $x = x \cdot Prob$

Partition fn for particles in magnetic field

$$Z = e^{-\beta \vec{E}} = e^{-\beta(\vec{\mu}\cdot\vec{H})} = e^{-\beta(\vec{\mu}\cdot\vec{B})}$$

up → μB, down → $-\mu B$

$$Z = e^{-\mu\beta B} + e^{\mu\beta B}$$

EXAMPLE → If $F = NkT \ln\left(\frac{a_0 V}{N}(kT)^{5/2}\right)$ find U.

Soln. As $F = -NkT \ln z$

$$\Rightarrow z = \frac{-a_0 V}{N}(kT)^{5/2}$$

$$U = -\frac{\partial}{\partial\beta}(\ln z) = -\frac{\partial}{\partial\beta}\left(\ln\left(\frac{-a_0 V}{N}\left(\frac{1}{\beta}\right)^{5/2}\right)\right)$$

$$= \frac{5N}{2\beta} = \frac{5}{2}NkT$$

$$\boxed{U = \frac{5}{2}NkT}$$

Ensemble → system of collection of similar systems having some const parameter

Microcanonical	Canonical	Grand canonical
Walls are rigid impermeable well insulated	Wallsar rigid conducting impenetrable	Walls are impenetrable rigid conducting

N,V,E		V,N,T		P,V,T
N,V,E		V,N,T		P,V,T
N,V,E		V,N,T		P,V,T
N,V,E				
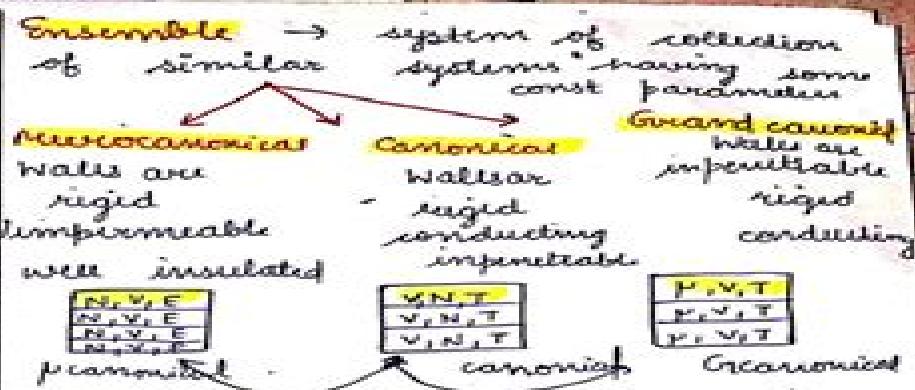
μcanonical   canonical   Grand canonical

These quantities remain same in a particle ensemble i·e in micro canonical, canonical & Grand canonical resp $\frac{1}{\sqrt{N}}$ ← Energy fluctuation → $\frac{1}{N}$

Maxwell Boltzmann Dist^n law of velocities

No. of particles in velocity range b/w v_x & $v_x + dx$ is

$$n(v_x)dv_x = N\left(\frac{m}{2\pi kT}\right)^{1/2} e^{-\frac{mv_x^2}{2kT}}dv_x$$

$$P = \frac{n}{N} = \left(\frac{m}{2\pi kT}\right)^{1/2} e^{-\frac{mv_x^2}{2kT}}dv_x$$

$$P_{max} = \left(\frac{m}{2\pi kT}\right)^{1/2} \quad \boxed{P \propto T^{-1/2}}$$

$$\bar{v_x} = \int_{-\infty}^{\infty} v_x P \, dv_x = 0$$

$$\boxed{\bar{v_x} = \bar{v_y} = \bar{v_z} = 0}$$

$$\boxed{\langle v_x v_y\rangle = \langle v_y v_z\rangle = \langle v_z v_x\rangle = 0}$$

In 1-D

$$\langle v_x^2\rangle = \langle v_y^2\rangle = \langle v_z^2\rangle = kT/m$$

In 2-D

$$\langle v_x^2\rangle = \langle v_y^2\rangle = \langle v_z^2\rangle = 2kT/m$$

In 3-D

$$\langle v_x^2\rangle = \langle v_y^2\rangle = \langle v_z^2\rangle = 3kT/m$$

Statistical Physics

L D Landau,E.M. Lifshitz

Statistical Physics:

Basics Of Statistical Physics: A Bachelor Degree Introduction Harald J W Muller-kirsten,2009-08-21 Statistics links microscopic and macroscopic phenomena and requires for this reason a large number of microscopic elements like atoms The results are values of maximum probability or of averaging This introduction to statistical physics concentrates on the basic principles and attempts to explain these in simple terms supplemented by numerous examples The basic principles concentrated on are the difference between classical and quantum statistics the a priori probabilities as related to degeneracies the vital aspect of indistinguishability as compared with distinguishability in classical physics the differences between conserved and nonconserved elements the latter including photons and phonons the different ways of counting arrangements in the three statistics Maxwell Boltzmann Fermi Dirac Bose Einstein the difference between maximization of the number of arrangements of elements in these and averaging in the Darwin Fowler method Significant applications to solids radiation and to electrons in metals are treated in separate chapters Finally the Bose Einstein distribution is rederived under condensation conditions Each chapter concludes with examples and exercises *Statistical Physics* Franz Mandl,1991-01-08 The Manchester Physics Series General Editors D J Sandiford F Mandl A C Phillips Department of Physics and Astronomy University of Manchester Properties of Matter B H Flowers and E Mendoza Optics Second Edition F G Smith and J H Thomson Statistical Physics Second Edition E Mandl Electromagnetism Second Edition I S Grant and W R Phillips Statistics R J Barlow Solid State Physics Second Edition J R Hook and H E Hall Quantum Mechanics F Mandl Particle Physics Second Edition B R Martin and G Shaw The Physics of Stars Second Edition A C Phillips Computing for Scientists R J Barlow and A R Barnett Statistical Physics Second Edition develops a unified treatment of statistical mechanics and thermodynamics which emphasises the statistical nature of the laws of thermodynamics and the atomic nature of matter Prominence is given to the Gibbs distribution leading to a simple treatment of quantum statistics and of chemical reactions Undergraduate students of physics and related sciences will find this a stimulating account of the basic physics and its applications Only an elementary knowledge of kinetic theory and atomic physics as well as the rudiments of quantum theory are presupposed for an understanding of this book Statistical Physics Second Edition features A fully integrated treatment of thermodynamics and statistical mechanics A flow diagram allowing topics to be studied in different orders or omitted altogether Optional starred and highlighted sections containing more advanced and specialised material for the more ambitious reader Sets of problems at the end of each chapter to help student understanding Hints for solving the problems are given in an Appendix

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